Angle deformation of Kähler–Einstein edge metrics on Hirzebruch surfaces

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Dedicated to Bernie Shiffman on the occassion of his retirement

Abstract

We construct a family of Kähler–Einstein edge metrics on all Hirzebruch surfaces using the Calabi ansatz and study their angle deformation. This allows us to verify in some special cases a conjecture of Cheltsov–Rubinstein that predicts convergence towards a non-compact Calabi–Yau fibration in the small angle limit. We also give an example of a Kähler–Einstein edge metric whose edge singularity is rigid, answering a question posed by Cheltsov.

1 Introduction

Recently, Cheltsov–Rubinstein [4] put forward a conjectural picture in which non-compact Calabi–Yau fibrations emerge as the small angle limit of families of compact singular metrics known as Kähler edge metrics. In a recent article, we verfied this picture in the most elementary setting of Riemann surfaces [12]. Our goal in the present article is to give further evidence for the conjecture by verifying it for some special symmetric complex surfaces. In passing, we also answer a question of Cheltsov by proving the existence of a rigid Kähler–Einstein edge metric.

Let M be a smooth complex projective manifold and let $D = D_1 + \ldots + D_r$ be a simple normal crossing divisor in M. Following Tian [14, p. 147], a Kähler metric on M with an edge singularity along a divisor $D \subset M$ is a smooth Kähler metric on $M \setminus D$ that has a cone singularity transverse to D ('bent' at an angle $2\pi\beta_i < 2\pi$ along D_i). Such metrics are called *Kähler edge metrics;* we refer to [9] for geometric analysis on such spaces and to [11] for a detailled survey and further references.

In op. cit. [4] one is concerned with small cone angle Kähler edge metrics that are also Einstein, called Kähler–Einstein edge (KEE) metrics and of positive Ricci curvature. For such metrics to exist it is necessary that (X, D) be *asymptotically log Fano* [4, Definition 1.1], a positivity property generalizing positivity of the first Chern class (i.e., being Fano) in the sense that $c_1(M) - \sum_{i=1}^{r}(1 - \beta_i(j))D_i$ is positive for a sequence $\beta(j) = (\beta_1(j), \ldots, \beta_r(j)) \in (0,1)^r$ tending to the origin (being Fano corresponds to the other 'extreme' with $\beta_1 = \ldots = \beta_r = 1$ or alternatively there being no divisor, i.e., $D = 0$). Assuming then that (M, D) is asymptotically log Fano and that it admits KEE metrics $\{\omega_{\beta(j)}\}_{j \in \mathbb{N}}$ with angles $\beta_i(j)$ along D_i, it is conjectured there that by taking an appropriate limit these metrics will converge to a Calabi–Yau fibration on the non-compact space $M \setminus D$ [4, Conjecture 1.11].

Conjecturally, the fibration may be described in terms of the adjoint anticanonical linear system as follows. Let K_M denote the canonical bundle, so that $c_1(M) = -c_1(K_M)$ by definition. By our assumption $-K_M - D$ is nef as it is a limit of ample divisors. Let

$d = \dim H^0(M, -K_M - D)$ and let $s_1, \ldots, s_d$ be a basis for the vector space $H^0(M, -K_M - D)$. The Kodaira map

$$M \ni z \mapsto [s_1(z) : \ldots : s_d(z)] \in \mathbb{P}^{d-1}, \tag{1.1}$$

is then a holomorphic map from M onto its image, a projective variety we denote by $Y \subset \mathbb{P}^{d-1}$. Denote by $\kappa := \dim Y$. It was conjectured [4, Conjecture 1.6] that $\kappa < \dim M$ (under the assumption that the KEE metrics ω_β exist for small β), or equivalently that $(K_M + D)^{\dim M} = 0$, and this was established in dimension 2 [5] and subsequently in all dimensions [6]. The next step in the program described in [4] is to study the geometric limit, if such exists, of such KEE metrics when the angles tend to zero. More precisely, we are interested in the following conjecture [4, Conjecture 1.8] (as customary, when discussing dimension 2 we will use 'del Pezzo' instead of 'Fano', and replace (M, D) by (S, C)):

Conjecture 1.1. *Let $C = C_1 + \ldots + C_r \not\sim -K_S$ be a disjoint union of smooth curves in a projective surface S such that (S, C) is strongly asymptotically log del Pezzo. Let $\beta \in (0, 1)^r$ be a sequence tending to the origin and suppose that there exist Kähler–Einstein edge metrics ω_β of angle $2\pi\beta_i$ along D_i and of positive Ricci curvature. Then as $\beta \to 0$ an appropriate limit of ω_β converges to a fibration of cylinders whose base is $\mathbb{P}^1$.*

In fact, by general results of Kawamata and Shokurov C can have at most two connected components, so that $r \in \{1, 2\}$ [4, Remark 3.7]. Note that part of the challenge is to identify what exactly "appropriate" means in Conjecture 1.1. As we show in this note the appropriate sense, at least in some cases, turns out to be rescaling the fibers of the fibration (1.1) in a precise way and considering a certain pointed Gromov–Haussdorff limit, inspired by the asymptotic analysis in the Riemann surface case [12]. Moreover, we also show that the *un-rescaled* limit exists as a collapsed Gromov–Haussdorff limit, and is the Kähler–Einstein metric on the base Y (in this case the Fubini–Study metric on $\mathbb{P}^1$). We believe this should generalize to the setting of Conjecture 1.1.

In this note we concentrate on a particular family of (S, C) as in Conjecture 1.1 and first construct rather explicitly the sequence of KEE metrics ω_β for a sequence of small angles $\beta \in (0, 1)^2$ tending to zero. Moreover, we analyze the small angle limit of this sequence and resolve Conjecture 1.1 for this family of pairs. The pairs we consider are

$$S = \mathbb{F}_n, \, n \in \mathbb{N} \cup \{0\}, \quad C_1 = Z_n, \, C_2 \in |Z_n + nF|, \tag{1.2}$$

where $\mathbb{F}_n$ is the n-th Hirzebruch surface [7], C_1 is the $-n$-curve, and C_2 is an n-curve; see §2.3 and §3.3 for notation and precise definitions. The 'boundary' in these pairs consists of two disjoint components C_1 and C_2 which in some senses is an added difficulty, however, the pairs (1.2) have the advantage of being toric (by which we mean that S is toric and each C_i is a torus-invariant holomorphic curve in S), and in fact the pairs even admit the larger symmetry group of the Calabi ansatz. The simplest sub-case, $n = 0$, serves as a general guide, and, in fact, as a nice bridge to our previous 1-dimensional work [12]. In the $n = 0$ case, the boundary consists of two disjoint fibers in $\mathbb{P}^1 \times \mathbb{P}^1$ and while the class $-K_{\mathbb{P}^1 \times \mathbb{P}^1} - (1 - \beta_1)F_1 - (1 - \beta_2)F_2$ is ample for all small β_1 and β_2, only the classes with $\beta_1 = \beta_2$ admit KEE metrics which we denote by $\omega_{\beta_1, \beta_1}$. This is reminiscent of the situation for footballs, and for a good reason. In fact the reason that only the classes with $\beta_1 = \beta_2$ admit KEE metrics is that if $\beta_1 \neq \beta_2$ then $\mathbb{P}^1$ admits a conic Ricci soliton football metric with angles $2\pi\beta_1$ at the North pole and $2\pi\beta_2$ at the South pole (see, e.g., [12, §2]). Taking the product with the Fubini–Study metric on $\mathbb{P}^1$ we see that $(\mathbb{P}^1 \times \mathbb{P}^1, Z_0 + Z_0)$ admits a Kähler–Ricci soliton edge metric with nontrivial vector field, hence cannot admit a KEE metric. When $\beta_1 = \beta_2$, the KEE metric $\omega_{\beta_1, \beta_1}$ is nothing

but the product of the volume one Fubini–Study metric on the base $Y = \mathbb{P}^1$ and the volume $4\pi\beta_1$ constant scalar curvature conic metric, i.e., the football, on the fibers. Then, our previous work [12, Theorem 1.3] shows that after rescaling the fiber metrics by $1/\beta_1^2$ there is a pointed Gromov–Haussdorff limit that converges to the product space $(\mathbb{P}^1 \times \mathbb{C}^\star, \pi_1^*\omega_{\text{FS}} + C\pi_2^*\omega_{\text{Cyl}})$ for some $C > 0$. In fact, this can also be shown by elementary complex analysis (op. cit. is concerned with the more general setting of Ricci solitons).

When $n > 0$ the results of [12] do not apply as the situation is no longer a product one and we need different tools. First, we describe some of the well-known Calabi ansatz computations that reduce the KEE equation to an ODE. These are classical when there are no edges, and as we show using the asymptotic expansion of [9, Theorem 1] they generalize naturally to the setting of edges. Similarly to the case $n = 0$ it turns out there is a relation between β_1 and β_2 that is necessary and sufficient for the class $-K_{\mathbb{F}_n} - (1 - \beta_1)C_1 - (1 - \beta_2)C_2$ to admit a KEE representative ω_{β_1,β_2}. Moreover, $\beta_2 \to 0$ as $\beta_1 \to 0$ and moreover $\lim_{\beta_1 \to 0} \beta_2/\beta_1 = 1$. We then generalize the asymptotic analysis from $n = 0$ to prove Conjecture 1.1 for the pairs (1.2) by analyzing rather explicitly such sequences of KEE metrics ω_{β_1,β_2} using the Calabi ansatz.

Theorem 1.2. *Consider the pairs (1.2). Then for each $n \in \mathbb{N}$:*

- *There exists a family of KEE metrics ω_{β_1,β_2} on the pairs (1.2) for each*

$$(\beta_1, \beta_2) = \Big(\beta_1, \frac{1}{2n}\big(n\beta_1 - 3 + \sqrt{3(3 - n\beta_1)(1 + n\beta_1)}\big)\Big), \quad \text{for all } \beta_1 \in (0, 2/n) \cap (0, 1],$$

with ω_{β_1,β_2} cohomologous to

$$\frac{2 + n\beta_2}{2 - n\beta_1}[C_2] - [C_1].$$

- *As β_1 tends to zero, $(\mathbb{F}_n, \omega_{\beta_1,\beta_2})$ converges in the Gromov–Haussdorff sense to $(\mathbb{P}^1, n\omega_{\text{FS}})$, where ω_{FS} is the Fubini–Study metric. The convergence also holds pointwise in all C^k norms on the level of tensors.*

- *The fiberwise-rescaled metrics $\widetilde{\omega_{\beta_1,\beta_2}}$ obtained by rescaling by β_1^{-2} only the restriction of ω_{β_1,β_2} to each fiber (of the projection to $\mathbb{P}^1$) converges in C^k on compact subsets as well as in the pointed Gromov–Haussdorff sense to a cylinder fibration over $\mathbb{P}^1$, $(\mathbb{P}^1 \times \mathbb{C}^\star, n\pi_1^*\omega_{\text{FS}} + n\pi_2^*\omega_{\text{Cyl}})$, i.e.,*

$$\lim_{\beta_1 \to 0} \widetilde{\omega_{\beta_1,\beta_2}} = n(\pi_1^*\omega_{\text{FS}} + \pi_2^*\omega_{\text{Cyl}}),$$

where ω_{Cyl} is the flat metric on $\mathbb{C}^\star$, and π_i is the projection on to the i-th factor.

In 2015, Cheltsov posed the following question [3].

Question 1.1. *Let D be a smooth connected divisor in a smooth variety X, and let $\beta \in (0, 1)$. Does there exists a triple (X, D, β) such that (X, D) admits a KEE metric of positive Ricci curvature and angle $2\pi\beta$ along D but does not admit KEE metrics of angles $2\pi\alpha$ along D for any $(0, 1) \ni \alpha \neq \beta$?*

Observe that the question becomes interesting in dimension two and above since there does not exist a constant scalar curvature conic Riemann sphere with a single cone point. We answer this question affirmatively.

Theorem 1.3. *Let Z_{-1} be a smooth curve of self-intersection 1 disjoint from the -1-curve in the first Hirzebruch surface $\mathbb{F}_1$. The pair $(\mathbb{F}_1, Z_{-1})$ admits a KEE metric of angle $2\pi\beta \in (0, 2\pi]$ along Z_{-1} if and only if $\beta = \sqrt{3} - 1$.*

1.1 Organization

In Section 2, we provide several useful viewpoints of Hirzebruch surfaces for the reader's convenience. In Section 3, we characterize Kähler edge metrics on Hirzebruch surfaces (Proposition 3.3) using the Calabi ansatz and the asymptotic analysis of [9]. In Section 4, we solve explicitly the Kähler–Einstein edge equation and determine the corresponding angles along the boundary divisors (Corollary 4.1), and prove Theorem 1.3. Finally, in Section 5, we study the small angle limits of the Kähler–Einstein edge metrics and prove Theorem 1.2.

2 Several descriptions of Hirzebruch surfaces

2.1 The $\mathbb{A}_n$-singularity

Let $n \in \mathbb{N}$. The simplest singularity in complex geometry is the $\mathbb{A}_n$ singularity constructed as follows. Let $\mathbb{Z}_n$ act on $\mathbb{C}^2$ by the diagonal action, so that the $\mathbb{Z}_n$-orbit of a point (a_1, a_2) is the collection of n points

$$\{(e^{2\pi\sqrt{-1}l/n}a_1, e^{2\pi\sqrt{-1}l/n}a_2) \in \mathbb{C}^2 : l = 0, \ldots, n-1\}.$$

This is always a collection of n distinct points, unless $(a_1, a_2) = (0, 0)$. So if we consider the orbit space (or quotient space)

$$\mathbb{C}^2/\mathbb{Z}_n$$

defined, as a set, as $\mathbb{C}^2$ quotiented by the equivalence relation

$$\begin{aligned} &(a_1, a_2) \sim (b_1, b_2) \text{ if there is some } l \in \{0, \ldots, n-1\} \\ &\quad \text{such that } a_1 = e^{2\pi\sqrt{-1}l/n}b_1 \text{ and } a_2 = e^{2\pi\sqrt{-1}l/n}b_2 \text{ (with the same } l\text{)}, \end{aligned} \tag{2.1}$$

we obtain an orbifold, smooth on the complement of a single singular point.

2.2 Blow-up description of the total space

In the remainder of this note we always assume

$$n \in \mathbb{N},$$

since the case $n = 0$ was treated in the Introduction. We resolve the $\mathbb{A}_n$ singularity by blowing up $\mathbb{C}^2/\mathbb{Z}_n$ at the single singular point. Next we give an alternative global description of the resulting resolution.

Denote by

$$\mathrm{Bl}_0(\mathbb{C}^2/\mathbb{Z}_n) \tag{2.2}$$

the blow-up at the origin of (the orbifold, if $n \geq 2$) $\mathbb{C}^2/\mathbb{Z}_n$. Denote by $H \to \mathbb{P}^1$ the hyperplane bundle over $\mathbb{P}^1$ and by $H^* \to \mathbb{P}^1$ the dual (tautological) line bundle. Then by

$$-nH_{\mathbb{P}^1} \tag{2.3}$$

we denote the (2-dimensional) total space of $-nH \equiv (H^*)^{\otimes n}$ considered as a line bundle over $\mathbb{P}^1$. The following is an elementary exercise.

Lemma 2.1. $\mathrm{Bl}_0(\mathbb{C}^2/\mathbb{Z}_n)$ *is biholomorphic to* $-nH_{\mathbb{P}^1}$.

2.3 Compactifying at infinity

By adding a point at infinity in each $\mathbb{C}$ fiber (in the $-kH_{\mathbb{P}^1}$ description) we obtain a compact space, a $\mathbb{P}^1$ fibration over $\mathbb{P}^1$. The space is equivalently obtained by taking a product of each $\mathbb{C}$ fiber with another copy of $\mathbb{C}$ and then taking the quotient under the $\mathbb{C}^*$ action on the $\mathbb{C}^2$ fibers (but not acting on the base). This space is often denoted by

$$\mathbb{P}\big(-nH_{\mathbb{P}^1} \oplus \mathbb{C}_{\mathbb{P}^1}\big), \tag{2.4}$$

and we will simply denote it by

$$\mathbb{F}_n, \tag{2.5}$$

for any $n \geq 0$. In effect, we have added a copy of $\mathbb{P}^1$. We therefore have two distinguished divisors in the space: the zero section

$$Z_n := \{w = 0\} \tag{2.6}$$

(the exceptional divisor in the blow-up description) and the infinity section

$$Z_{-n} := \{w = \infty\}. \tag{2.7}$$

By construction the two do not intersect.

We can take (2.4) as the definition of $\mathbb{F}_n$. In that case it actually makes sense for any $n \in \mathbb{Z}$. And, since for any vector bundle A and line bundle L we have $\mathbb{P}(A \oplus L) = \mathbb{P}(A)$, it follows that $\mathbb{F}_{-n}$ is biholomorphic to $\mathbb{F}_n$ (take $L = 2nH_{\mathbb{P}^1}$) with the biholomorphism exchanging the zero and the infinity sections (i.e., Z_n with Z_{-n}). Recalling Lemma 2.1 we have shown the following:

Lemma 2.2. $\mathbb{F}_n \setminus Z_{-n}$ is biholomorphic to $\mathrm{Bl}_0(\mathbb{C}^2/\mathbb{Z}_n)$, which is in turn biholomorphic to $\mathbb{F}_n \setminus Z_n$.

Lemma 2.3. $\dim H_{\bar{\partial}}^{1,1}(\mathbb{F}_n) = 2$.

Proof. By [7],

$$\mathbb{F}_n \xrightarrow[\text{diffeo.}]{\cong} \begin{cases} S^2 \times S^2, & n \text{ even,} \\ \mathbb{P}^2 \# \overline{\mathbb{P}^2}, & n \text{ odd.} \end{cases}$$

So one has $\dim H^2(\mathbb{F}_n, \mathbb{C}) = 2$. Thus $2\dim H_{\bar{\partial}}^{2,0}(\mathbb{F}_n) + \dim H_{\bar{\partial}}^{1,1}(\mathbb{F}_n) = 2$, so it follows that $\dim H_{\bar{\partial}}^{1,1}(\mathbb{F}_n)$ is either 0 or 2. The former possibility is excluded since $Z_n \cong \mathbb{P}^1$ is a non-trivial holomorphic $(1,1)$-cycle. $\square$

Remark 2.4. We will explicitly construct $(1,1)$-forms in Section 3.3 below which gives an alternative direct description why $\dim H_{\bar{\partial}}^{1,1}(\mathbb{F}_n, \mathbb{Z}) \geq 2$ hence equal to 2. In fact, we will construct explicit representatives for each Kähler class.

3 Edge metrics on Hirzebruch surfaces

3.1 Edge metrics on powers of the tautological line bundle

In this section we elaborate on Remark 2.4 and give a geometric approach to compactification.

We start by constructing some $U(2)$-invariant Kähler edge metrics on $-nH_{\mathbb{P}^1}$ (2.3). The construction goes back to Calabi [1, 2] who considered the case of smooth Kähler metrics. The generalization to the edge case is not much harder.

Consider the Hermitian metric h on $-H$ that assigns to each point in the total space (Z, w) the norm squared

$$|(Z, w)|_h^2 := |w|^2 ||Z||^2 = |w|^2(|Z_1|^2 + |Z_2|^2),$$

i.e., locally

$$h(Z) = ||Z||^2,$$

where $||Z||$ is the Euclidean norm of the vector (Z_1, Z_2) in $\mathbb{C}^2$. Similarly, h^n is a metric on $-nH$ and

$$|(Z, w)|_{h^n}^2 := |w|^2 ||Z||^2 = |w|^2(|Z_1|^2 + |Z_2|^2)^n,$$

and

$$h^k(Z) = (|Z_1|^2 + |Z_2|^2)^k.$$

On the chart $Z_2 \neq 0$, we choose local holomorphic coordinates $z := Z_1/Z_2$, so that

$$Z = [Z_1 : Z_2] = [z : 1]. \tag{3.1}$$

The curvature of this metric is a Kähler form on the base (i.e., the zero section $Z_n \cong \mathbb{P}^1$) which is given by

$$-\sqrt{-1}\partial\bar{\partial} \log h^n = -n\sqrt{-1}\partial\bar{\partial} \log ||Z||^2 = -n\sqrt{-1}\partial\bar{\partial} \, \log(1 + |z|^2)^2$$
$$= -n(1 + |z|^2)^{-2}\sqrt{-1}dz \wedge \overline{dz} = -n\omega_{\mathrm{FS}}.$$

In the sequel, we will use the above curvature property of h^n, as well as the fact that it is a globally defined $U(2)$-invariant smooth function on $\mathrm{Bl}_0(\mathbb{C}^2/\mathbb{Z}_n)$.

3.2   The Calabi ansatz on the total space

We use the logarithm of the global invariant function from above as our coordinate from now on. That is, we set

$$s(Z, w) := \log |(Z, w)|_{h^n}^2 = \log |w|^2 + n\log(1 + |z|^2), \quad (Z, w) \in -nH_{\mathbb{P}^1}, \tag{3.2}$$

and seek canonical Kähler metrics on $\mathrm{Bl}_0(\mathbb{C}^2/\mathbb{Z}_n)$ that depend solely on s, namely, Kähler metrics of the form

$$\eta = \sqrt{-1}\partial\bar{\partial}f(s),$$

where f is a smooth function. Our goal will be to determine appropriate f that make η have various desirable curvature properties and edge type singularities along $Z_{\pm n}$.

Denote

$$f'(s) := \frac{df}{ds}.$$

Note,

$$\frac{\partial s}{\partial w} = \frac{1}{w}, \quad \frac{\partial s}{\partial Z_i} = \frac{\overline{Z_i}}{||Z||^2}, \quad i = 1, 2.$$

Working on the chart $Z_2 = 1$ (recall (3.1)),

$$\begin{pmatrix} \dfrac{\partial^2 f}{\partial w \partial \bar{w}} & \dfrac{\partial^2 f}{\partial w \partial \bar{z}} \\ \dfrac{\partial^2 f}{\partial z \partial \bar{w}} & \dfrac{\partial^2 f}{\partial z \partial \bar{z}} \end{pmatrix} = \begin{pmatrix} \dfrac{f''}{|w|^2} & \dfrac{nf''z}{w(|z|^2 + 1)} \\ \dfrac{nf''\bar{z}}{\overline{w}(|z|^2 + 1)} & \dfrac{nf' + n^2f''|z|^2}{(|z|^2 + 1)^2} \end{pmatrix}$$

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Set
$$\pi_1(Z, w) := z, \quad \pi_2(Z, w) := w,$$
and denote by ω_{FS} the Fubini–Study metric on $\mathbb{P}^1$ and the flat Kähler form on the cylinder $\mathbb{C}^\star$ by
$$\omega_{\mathrm{Cyl}} := \frac{\sqrt{-1}dw \wedge d\bar{w}}{|w|^2}. \tag{3.3}$$
Thus,
$$\begin{aligned}
\eta = {} & nf'\pi_1^*\omega_{\mathrm{FS}} + f''\pi_2^*\omega_{\mathrm{Cyl}} \\
& + n\frac{f''}{1+|z|^2}\Big(\sqrt{-1}\frac{dw}{w} \wedge z\,\overline{dz} + \sqrt{-1}\bar{z}dz \wedge \frac{\overline{dw}}{\bar{w}}\Big) \\
& + n^2 f''(s)(1+|z|^2)^{-2}\bar{z}dz \wedge z\,\overline{dz}.
\end{aligned} \tag{3.4}$$
From this computation we see that the two key quantities are f' and f'' (rather than f itself). Both of these must be positive, i.e., f must be an increasing convex function of s. Inspired by this, consider a Legendre type change of variables going back to Calabi [1] (cf. [8]; see [12] for a reference most closely following our notation):
$$\tau = \tau(s) := f'(s), \quad \varphi = \varphi(\tau) = \varphi(\tau(s)) := f''(s), \quad \tau \in \mathrm{Im}f' = f'(\mathbb{R}). \tag{3.5}$$
Setting
$$\alpha := n\frac{\bar{z}dz}{1+|z|^2}, \tag{3.6}$$
and using (3.5), we may rewrite (3.4) as
$$\eta = n\tau\pi_1^*\omega_{\mathrm{FS}} + \varphi\Big(\pi_2^*\omega_{\mathrm{Cyl}} + \sqrt{-1}\alpha \wedge \bar{\alpha} + \sqrt{-1}\alpha \wedge \overline{dw/w} + \sqrt{-1}dw/w \wedge \bar{\alpha}\Big). \tag{3.7}$$

3.3 Kähler classes

By construction, the normal bundle of $Z_{\pm n}$ is $\mathcal{O}_{\mathbb{P}^1}(\mp n) = \mp nH_{\mathbb{P}^1}$, so
$$Z_{\pm n}^2 = \mp n, \quad Z_n.Z_{-n} = 0.$$
It follows from this and Lemma 2.3 that Z_n and Z_{-n} generate the Picard group, and thus by the Nakai–Moishezon criterion the Kähler classes are precisely represented by $-xZ_n + yZ_{-n}$ with $y > x > 0$ (this also follows directly from the Calabi ansatz). The divisor class
$$F := \frac{1}{n}\big(Z_{-n} - Z_n\big), \tag{3.8}$$
has zero self intersection and intersects $Z_{\pm n}$ exactly at one point, hence represents the fibers of the projection to the base Z_n. Setting
$$C_1 = Z_n, \quad C_2 = Z_{-n}, \tag{3.9}$$
and using (3.8) we recover the notation (1.2). The canonical class can be determined as follows. Write $K_{\mathbb{F}_n} = -xZ_n + yZ_{-n}$. By Riemann–Roch,
$$-2 = (K_{\mathbb{F}_n}+Z_{\pm n}).Z_{\pm n} = ((1-x)Z_n+yZ_{-n}).Z_n = -n(1-x) = (Z_n+(1+y)Z_{-n}).Z_{-n} = n(1+y),$$
so $x = 1 - 2/n$, $y = -1 - 2/n$, and
$$-K_{\mathbb{F}_n} \sim \Big(1-\frac{2}{n}\Big)Z_n + \Big(1+\frac{2}{n}\Big)Z_{-n} \sim 2Z_n + (n+2)F. \tag{3.10}$$

3.4 The angle constraint as boundary data

We rewrite the angle constraint in terms of τ. Note that the domain of τ is

$$\mathrm{Im}\, f' = (\inf f', \sup f'),$$

and τ must be positive so $\inf f' \geq 0$.

Lemma 3.1. *Suppose that η restricts to a Riemannian metric on E. Then, $\inf f' > 0$.*

Proof. If $\inf f' = 0$, then by (3.4) and (2.6) the restriction of η to E is identically zero, which means the zero section is collapsed to a point, a contradiction. $\square$

To simplify computations we will henceforth assume we are in the situation of Lemma 3.1 and rescale η, equivalently f, by a positive constant so that

$$\inf f' = 1. \tag{3.11}$$

This is equivalent to rescaling the Kähler class of η. Thus the only contribution to the Poincaré–Lelong formula aside from ω_{Cyl} will be from any vanishing of φ along $\{w = 0\}$ and along $\{1/w = 0\}$.

By construction, $Z_{\pm n} = \{s = \mp\infty\}$. By our normalization above τ ranges in the domain $(1, T)$ with

$$Z_n = \{\tau = 1\} \quad Z_{-n} = \{\tau = T\}. \tag{3.12}$$

Lemma 3.2. $T < \infty$.

Proof. By assumption η is a Kähler edge metric on the pair (1.2). The distance function of η depends only on s. Restricting η to a fiber (i.e., say, to the vertical section $\{z = 0\}$) using (3.7), we get an S^1-invariant metric

$$g = \frac{1}{2\varphi(\tau)} d\tau^2 + 2\varphi(\tau) d\theta^2, \tag{3.13}$$

where $w = e^{s/2 + \sqrt{-1}\theta}$ is a coordinate on $\{z = 0\}$. This follows in the same way as in the 1-dimensional setting [12, Lemma 2.1]. Note that here we implicitly used the fact that η is $U(2)$-invariant, hence its restriction to any fiber is S^1-invariant. Since for any Kähler edge metric the distance to the boundary is finite [9, Remark 6.4], we see from (3.12) and (3.13) that Z_{-n} is at finite distance from Z_n if and only if $\int_1^T \sqrt{\varphi(\tau)} d\tau < \infty$. If $T = \infty$ by (3.5) it follows that also $\lim_{s \to \infty} f''(s) = \infty$ so $\lim_{\tau \to T} \varphi(\tau) = \infty$. It follows that τ ranges in a finite interval. $\square$

By our assumption (3.11) and Lemma 3.2, τ ranges in a domain $(1, T)$ with $T < \infty$. Thus,

$$\lim_{s \to \pm\infty} \frac{d\tau}{ds} = 0,$$

i.e., using (3.5),

$$\varphi(1) = \varphi(T) = 0. \tag{3.14}$$

Next, we rewrite the angle constraint at the edges in terms of τ.

Proposition 3.3. *η is a Kähler edge metric on the pair (1.2) with angle $2\pi\beta_1$ along C_1 and $2\pi\beta_2$ along C_2 if and only if*

$$\varphi(1) = 0, \quad \varphi'(1) = \beta_1, \quad \varphi(T) = 0, \quad \varphi'(T) = -\beta_2. \tag{3.15}$$

Proof. By (3.14) it remains to determine the derivatives of φ at 1 and T.

Suppose first that η is a Kähler edge metric with angles as stated. It follows from [9, Theorem 1, Proposition 4.4] that f has complete asymptotic expansions both near $w = 0$ and $w = \infty$. Let us concentrate on the former first. The leading term in that expansion is $|w|^{2\beta_1}$ and using (3.2),

$$\varphi \sim C_1 + C_2|w|^{2\beta_1} + (C_3 \sin\theta + C_4 \cos\theta)|w|^2 + O(|w|^{2+\epsilon})$$
$$= C_1 + C_2 e^{\beta_1 s} + (C_3 \sin\theta + C_4 \cos\theta)e^s + O(e^{(1+\epsilon)s})$$

(note that r in [9, (56)] is equal to $|w|^{\beta_1}/\beta_1$ in our notation, see [9, p. 102]). Note that $C_1 = 0$ by (3.14) (actually also $C_3 = C_4 = 0$ as φ is independent of θ but we do not need this). Moreover, the expansion can be differentiated term-by-term as $|w| \to 0$ or $s \to -\infty$. As $\varphi'(\tau) = \frac{\partial\varphi}{\partial s}\frac{ds}{d\tau} = \frac{\partial\varphi}{\partial s}/\varphi$, we obtain

$$\varphi(1) = 0, \quad \varphi'(1) = \beta_1. \tag{3.16}$$

The same arguments imply that

$$\varphi(T) = 0, \quad \varphi'(T) = -\beta_2, \tag{3.17}$$

the minus sign coming from the fact that the leading term in the expansion is now $1/|w|^{2\beta_2} = e^{-\beta_2 s}$.

Conversely, suppose that (3.15) holds. Then near $Z_n = \{\tau = 1\} = \{w = 0\}$ there exists a positive smooth function $F(z,w)$ and a positive constant δ so that

$$\varphi(\tau(z,w)) = |w|^{2\delta}F(z,w). \tag{3.18}$$

Using (3.2), $|w|^{2\delta} = e^{\delta s}/(1+|z|^2)^\delta$, and

$$\frac{d}{d\tau} = \frac{1}{\varphi}\frac{d}{ds}, \tag{3.19}$$

so in the notation of (3.18),

$$\varphi_\tau(1) = \lim_{s \to -\infty} \frac{\frac{d}{ds}\left(F(s)e^{\beta s}\right)}{F(s)e^{\beta s}}$$
$$= \lim_{s \to -\infty} \frac{F'(s)e^{\beta s} + \beta F(s)e^{\beta s}}{F(s)e^{\beta s}}$$
$$= \beta + \lim_{s \to -\infty} F'(s)/F(s).$$

By assumption, $\lim_{s\to-\infty} F(s)$ is some positive (finite) number, in particular we must have $\lim_{s\to-\infty} F'(s) = 0$ Thus, $d\varphi/d\tau(1) = \delta$ and

$$\delta = \beta_1. \tag{3.20}$$

Now, taking the top wedge product of (3.7) gives,

$$\eta^2 = 2\left(n\tau\varphi(\tau)\pi_1^*\omega_{\mathrm{FS}} \wedge \pi_2^*\omega_{\mathrm{Cyl}} + \varphi^2\left(\pi_2^*\omega_{\mathrm{Cyl}} \wedge \sqrt{-1}\alpha \wedge \overline{\alpha} + \sqrt{-1}\alpha \wedge \overline{dw/w} \wedge \sqrt{-1}dw/w \wedge \overline{\alpha}\right)\right)$$
$$= 2n\tau\varphi(\tau)\pi_1^*\omega_{\mathrm{FS}} \wedge \pi_2^*\omega_{\mathrm{Cyl}}. \tag{3.21}$$

Now (3.3), (3.18) and (3.20) imply that η satisfies a complex Monge–Ampère equation with right-hand side equal to a smooth volume form times $1/|w|^{2-2\beta_1}$ near Z_n. Thus, [9, Theorem 1] applies and η is a Kähler edge metric with a complete asymptotic expansion near Z_n (note that op. cit. is stated in the case of a smooth connected divisor but applies verbatim in the case of smooth disjoint divisors). The same arguments apply near Z_{-n} to conclude. $\square$

4 The Einstein constraint

From the proof of Proposition 3.21 we can derive a formula for the Ricci tensor of η. Indeed, by (3.21) and the Poincaré–Leong formula,

$$\operatorname{Ric} \eta = (1 - \beta_1)[C_1] + (1 - \beta_2)[C_2] + 2\pi_1^* \omega_{\mathrm{FS}} - \sqrt{-1}\partial\bar{\partial} \log \tau - \sqrt{-1}\partial\bar{\partial} \log \varphi, \tag{4.1}$$

where the last term is understood to be the restriction of $\sqrt{-1}\partial\bar{\partial} \log \varphi$ to the complement of $C_1 + C_2$. For the remaining terms, compute

$$
\begin{aligned}
\partial_w \log f' &= \frac{f''}{f'}\frac{\partial s}{\partial w} = \frac{f''}{w f'} \\
\partial_z \log f' &= \frac{f''}{f'}\frac{\partial s}{\partial z} = n\frac{f''}{f'}\frac{\bar{z}}{1 + |z|^2} \\
\partial_{w\bar{w}} \log f' &= \left(\frac{f''}{f'}\right)'\frac{1}{w}\frac{\partial s}{\partial\bar{w}} = \left(\frac{f''}{f'}\right)'\frac{1}{|w|^2} \\
\partial_{z\bar{w}} \log f' &= n\left(\frac{f''}{f'}\right)'\frac{\overline{z/w}}{1 + |z|^2} \\
\partial_{z\bar{z}} \log f' &= n^2\left(\frac{f''}{f'}\right)'\frac{|z|^2}{(1 + |z|^2)^2} + n\frac{f''}{f'}\frac{1}{(1 + |z|^2)^2},
\end{aligned}
\tag{4.2}
$$

and this can be simplified by noting that

$$\frac{d\tau}{ds} = f'' = \varphi(\tau). \tag{4.3}$$

Thus,

$$
\begin{aligned}
\partial_w \log f' &= \frac{\varphi}{w\tau} \\
\partial_{z_i} \log f' &= n\frac{\varphi}{\tau}\frac{\bar{z}}{1 + |z|^2} \\
\partial_{w\bar{w}} \log f' &= \left(\frac{\varphi}{\tau}\right)_\tau\frac{\varphi}{|w|^2} \\
\partial_{z\bar{w}} \log f' &= n\left(\frac{\varphi}{\tau}\right)_\tau\varphi\frac{\overline{z/w}}{1 + |z|^2} \\
\partial_{|z|^2} \log f' &= n^2\left(\frac{\varphi}{\tau}\right)_\tau\varphi\frac{|z|^2}{(1 + |z|^2)^2} + n\frac{\varphi}{\tau}\frac{1}{(1 + |z|^2)^2},
\end{aligned}
\tag{4.4}
$$

Next, $\sqrt{-1}\partial\bar{\partial} \log \varphi$ is computed similarly by replacing f' by f'' everywhere. A simplification is obtained by noting that

$$\frac{f'''}{f''} = \frac{\varphi_\tau \varphi}{\varphi} = \varphi_\tau.$$

Thus,

$$\partial_w \log f'' = \frac{\varphi_\tau}{w}$$

$$\partial_z \log f'' = n\frac{\varphi_\tau \bar{z}}{1 + |z|^2}$$

$$\partial_{w\bar{w}} \log f'' = \frac{\varphi_{\tau\tau}\varphi}{|w|^2} \tag{4.5}$$

$$\partial_{z\bar{w}} \log f'' = n\frac{\varphi_{\tau\tau}\varphi\overline{z/w}}{1 + |z|^2}$$

$$\partial_{z\bar{z}} \log f'' = n^2\frac{\varphi_{\tau\tau}\varphi|z|^2}{(1 + |z|^2)^2} + n\varphi_\tau\frac{1}{(1 + |z|^2)^2}.$$

Altogether,

$$\sqrt{-1}\partial\bar\partial \log \tau = n\varphi/\tau\pi_1^*\omega_{\text{FS}}$$
$$+ \varphi(\varphi/\tau)_\tau\Big(\pi_2^*\omega_{\text{Cyl}} + \sqrt{-1}\alpha \wedge \bar\alpha \tag{4.6}$$
$$+ \sqrt{-1}\alpha \wedge \overline{dw/w} + \sqrt{-1}dw/w \wedge \bar\alpha.\Big),$$

and

$$\sqrt{-1}\partial\bar\partial \log \varphi = n\varphi_\tau\pi_1^*\omega_{\text{FS}}$$
$$+ \varphi\varphi_{\tau\tau}\Big(\pi_2^*\omega_{\text{Cyl}} + \sqrt{-1}\alpha \wedge \bar\alpha \tag{4.7}$$
$$+ \sqrt{-1}\alpha \wedge \overline{dw/w} + \sqrt{-1}dw/w \wedge \bar\alpha.\Big),$$

So, with α given by (3.6)

$$\text{Ric}\,\eta = (1 - \beta_1)[C_1] + (1 - \beta_2)[C_2] + (2 - n\varphi/\tau - n\varphi_\tau)\pi_1^*\omega_{\text{FS}}$$
$$- \varphi(\varphi/\tau + \varphi_\tau)_\tau\Big(\pi_2^*\omega_{\text{Cyl}} + \sqrt{-1}\alpha \wedge \bar\alpha \tag{4.8}$$
$$+ \sqrt{-1}\alpha \wedge \overline{dw/w} + \sqrt{-1}dw/w \wedge \bar\alpha\Big).$$

The Einstein edge equation

$$\text{Ric}\,\eta = \lambda\eta + (1 - \beta_1)[C_1] + (1 - \beta_2)[C_2] \tag{4.9}$$

becomes, using (3.7) and (4.8), the pair of equations

$$2 - n\varphi/\tau - n\varphi_\tau = n\lambda\tau$$
$$-\varphi\big(\varphi/\tau + \varphi_\tau\big)_\tau = \lambda\varphi.$$

Observe that the first equation implies the second by differentiating in τ. Also, setting $\tau = 1$ and using (3.15) implies

$$\lambda = \frac{2}{n} - \beta_1. \tag{4.10}$$

Observe that this already puts a constraint, as we must require positive Ricci curvature on $\mathbb{F}_n \setminus C$, equivalently

$$\beta_1 \in \left(0, \frac{2}{n}\right) \cap (0, 1]. \tag{4.11}$$

Thus, the Einstein equation near Z_n reduces to the first-order initial value problem

$$\varphi_\tau + \frac{\varphi}{\tau} = \frac{2}{n} + \left(\beta_1 - \frac{2}{n}\right)\tau, \qquad \varphi(1) = 0. \tag{4.12}$$

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4.1 Solving the Einstein equation

Using an integration factor τ this becomes

$$(\tau\varphi)_\tau = \frac{2}{n}\tau + \left(\beta_1 - \frac{2}{n}\right)\tau^2, \qquad \varphi(1) = 0, \tag{4.13}$$

so

$$\varphi = \frac{1}{n}\frac{\tau^2 - 1}{\tau} + \frac{1}{3}\left(\beta_1 - \frac{2}{n}\right)\frac{\tau^3 - 1}{\tau}. \tag{4.14}$$

Now, for this to correspond to a compact Kähler edge space with angle $2\pi\beta_2$ at Z_{-n} we must satisfy (3.15). To that end, let us determine $T, \varphi(T)$, and $\varphi_\tau(T)$ in (3.15) from (4.14). We factor φ as

$$\varphi(\tau) = \frac{1}{3}\left(\beta_1 - \frac{2}{n}\right)(\tau - 1)(\tau - \alpha_1)(\tau - \alpha_2)/\tau, \tag{4.15}$$

with $\alpha_1 \le \alpha_2$. Then,

$$\frac{1}{3}\left(\beta_1 - \frac{2}{n}\right)(\tau - \alpha_1)(\tau - \alpha_2) = \frac{1}{3}\left(\beta_1 - \frac{2}{n}\right)(\tau^2 + \tau + 1) + \frac{1}{n}(\tau + 1),$$

so

$$-\alpha_1\alpha_2 = \alpha_1 + \alpha_2 = \frac{1 + n\beta_1}{2 - n\beta_1}.$$

By (4.11) we see that $\alpha_1 < 0 < \alpha_2$, so $T = \alpha_2$ if we can show $\alpha_2 > 1$. Solving the quadratic equation for α_1, α_2 gives

$$\begin{aligned}
2\alpha_2 &= \frac{1 + n\beta_1 + \left[(1 + n\beta_1)^2 + 4(1 + n\beta_1)(2 - n\beta_1)\right]^{1/2}}{2 - n\beta_1} \\
&= \frac{1 + n\beta_1 + [(9 + 6n\beta_1 - 3n^2\beta_1^2]^{1/2}}{2 - n\beta_1}.
\end{aligned} \tag{4.16}$$

We claim that $\alpha_2 > 1$. According to (4.16), that amounts to verifying

$$\frac{1 + n\beta_1 + [(9 + 6n\beta_1 - 3n^2\beta_1^2]^{1/2}}{2 - n\beta_1} > 2,$$

which, after some manipulation precisely reduces to (4.11), proving the claim.

Finally, it remains to compute the final equation in (3.15), i.e., $\varphi_\tau(\alpha_2)$. Using (4.15),

$$\begin{aligned}
\varphi_\tau(\alpha_2) &= \frac{1}{3}\left(\beta_1 - \frac{2}{n}\right)(\tau - 1)(\tau - \alpha_1)/\tau\Big|_{\tau=\alpha_2} \\
&= \frac{1}{3}\left(\beta_1 - \frac{2}{n}\right)(\alpha_2 - 1)(\alpha_2 - \alpha_1)/\alpha_2.
\end{aligned} \tag{4.17}$$

Note that

$$\begin{aligned}
\alpha_2 - 1 &= \frac{\sqrt{3(3 - n\beta_1)(1 + n\beta_1)} + 3(n\beta_1 - 1)}{4 - 2n\beta_1}, \\
\frac{\alpha_2 - \alpha_1}{\alpha_2} &= \frac{2\sqrt{3(3 - n\beta_1)(1 + n\beta_1)}}{1 + n\beta_1 + \sqrt{3(3 - n\beta_1)(1 + n\beta_1)}},
\end{aligned} \tag{4.18}$$

Thus,

$$
\begin{aligned}
\varphi_\tau(\alpha_2) &= \frac{n\beta_1 - 2}{3n} \frac{\sqrt{3(3-n\beta_1)(1+n\beta_1)} + 3(n\beta_1 - 1)}{4 - 2n\beta_1} \frac{2\sqrt{3(3-n\beta_1)(1+n\beta_1)}}{1 + n\beta_1 + \sqrt{3(3-n\beta_1)(1+n\beta_1)}} \\
&= -\frac{1}{3n} \frac{\sqrt{3(3-n\beta_1)(1+n\beta_1)}\left(\sqrt{3(3-n\beta_1)(1+n\beta_1)} + 3(n\beta_1 - 1)\right)}{1 + n\beta_1 + \sqrt{3(3-n\beta_1)(1+n\beta_1)}} \\
&= -\frac{1}{3n} \frac{3(3-n\beta_1)(1+n\beta_1) + 3(n\beta_1 - 1)\sqrt{3(3-n\beta_1)(1+n\beta_1)}}{1 + n\beta_1 + \sqrt{3(3-n\beta_1)(1+n\beta_1)}} \\
&= -\frac{1}{3n}\left[\frac{[3(3-n\beta_1)(1+n\beta_1)]^{3/2} + 3(n\beta_1-1)3(3-n\beta_1)(1+n\beta_1)}{-(1+n\beta_1)^2 + 3(3-n\beta_1)(1+n\beta_1)} \right. \\
&\qquad\qquad \left. - \frac{3(3-n\beta_1)(1+n\beta_1)^2 + 3(n^2\beta_1^2-1)\sqrt{3(3-n\beta_1)(1+n\beta_1)}}{-(1+n\beta_1)^2 + 3(3-n\beta_1)(1+n\beta_1)} \right] \\
&= -\frac{1}{3n}\left[\frac{\sqrt{3(3-n\beta_1)(1+n\beta_1)}\left[3(3-n\beta_1)(1+n\beta_1) - 3(n^2\beta_1^2-1)\right]}{-4n^2\beta_1^2 + 8 + 4n\beta_1} \right. \\
&\qquad\qquad \left. + \frac{3(3-n\beta_1)(2n^2\beta_1^2 - 4 - 2n\beta_1)}{-4n^2\beta_1^2 + 8 + 4n\beta_1} \right] \\
&= \frac{3 - n\beta_1}{2n} + \frac{1}{n} \frac{\sqrt{3(3-n\beta_1)(1+n\beta_1)}\left[(3-n\beta_1)(1+n\beta_1) - (n^2\beta_1^2-1)\right]}{4n^2\beta_1^2 - 8 - 4n\beta_1} \\
&= \frac{3 - n\beta_1 - \sqrt{3(3-n\beta_1)(1+n\beta_1)}}{2n} = -\beta_2.
\end{aligned}
\tag{4.19}
$$

Note that $\beta_2 = \beta_2(\beta_1) < \beta_1$ for all β_1. Yet as β_1 tends to zero, β_2 tends to β_1, to wit,

$$
\begin{aligned}
\beta_2 &= \frac{n\beta_1 - 3 + 3\sqrt{1 + \frac{2}{3}n\beta_1 - \frac{1}{3}n^2\beta_1^2}}{2n} \\
&= \frac{1}{2n}\left(n\beta_1 - 3 + 3\left(1 + \frac{1}{3}n\beta_1 - \frac{1}{6}n^2\beta_1^2 - \frac{1}{8}\left(\frac{2}{3}n\beta_1 - \frac{1}{3}n^2\beta_1^2\right)^2 + O(\beta_1^3)\right)\right) \\
&= \beta_1 - \frac{n}{3}\beta_1^2 + O(\beta_1^3).
\end{aligned}
\tag{4.20}
$$

Combining all the above together with (3.10), (4.10), and (4.9) we have shown the following.

Corollary 4.1. *Let $n \in \mathbb{N}$. For each $\beta_1 \in (0, 2/n) \cap (0,1]$, there exists a Kähler–Einstein edge metric ω_{β_1,β_2} cohomologous to*

$$
\frac{2 + n\beta_2}{2 - n\beta_1}[Z_{-n}] - [Z_n] \sim \frac{n(\beta_1 + \beta_2)}{2 - n\beta_1}[Z_n] + n\frac{2 + n\beta_2}{2 - n\beta_1}[F]
$$

on the pair (1.2) with angles $2\pi\beta_1$ along $C_1 = Z_n$ and $2\pi\beta_2 = \pi\left(n\beta_1 - 3 + \sqrt{3(3-n\beta_1)(1+n\beta_1)}\right)/n$ along $C_2 = Z_{-n}$. One has $\beta_2 < \beta_1$ and $\lim_{\beta_1 \to 0} \beta_1/\beta_2 = 1$.

4.2 A rigid Kähler–Einstein edge metric

We now turn to proving Theorem 1.3.

In the case $n = 1$, β_1 can be taken *equal* to 1 by Corollary 4.1. This gives a KEE metric with angle $2\pi(\sqrt{3} - 1)$ along Z_{-1} and smooth elsewhere. This KEE metric *cannot* be deformed to small angles solely along Z_{-1}. In fact, $(\mathbb{F}_1, Z_{-1})$ is log del Pezzo (i.e., $-K_{\mathbb{F}_1} - Z_{-1}$ is ample)

so by [5, Proposition 3.1] it does not admit KEE metrics for *small* angles. But, in fact, more is true. By [5, Example 2.8] the pair $(\mathbb{F}_1, Z_{-1})$ does not admit KEE metrics for *any angle smaller* than $\sqrt{3}-1$ because Z_{-1} log slope destabilizes it, and neither for *any angle larger* than $\sqrt{3}-1$ because Z_1 log slope destabilizes it. This proves Theorem 1.3. This metric is quite remarkable in that it is in fact perhaps the first example of an "isolated" KEE metric whose edge singularity cannot be deformed at all.

5 Small angle limits

First, let us determine the un-rescaled small angle limit of the KEE metrics ω_{β_1,β_2}. By (4.18)

$$\alpha_2 = 1 + 3\frac{\sqrt{1 + \frac{2}{3}n\beta_1 - \frac{1}{3}n^2\beta_1^2} + n\beta_1 - 1}{4 - 2n\beta_1}$$
$$= 1 + 3\frac{\frac{4}{3}n\beta_1 - \frac{2}{9}n^2\beta_1^2 + O(\beta_1^3)}{4 - 2n\beta_1} = 1 + n\beta_1 + \frac{n^2}{3}\beta_1^2 + O(\beta_1^3), \tag{5.1}$$

while,

$$\alpha_1 = 1 + 3\frac{-\sqrt{1 + \frac{2}{3}n\beta_1 - \frac{1}{3}n^2\beta_1^2} + n\beta_1 - 1}{4 - 2n\beta_1}$$
$$= 1 + 3\frac{-2 + \frac{2}{3}n\beta_1 + \frac{2}{9}n^2\beta_1^2 + O(\beta_1^3)}{4 - 2n\beta_1} = -\frac{1}{2} + \frac{n}{2}\beta_1 + \frac{5n^2}{12}\beta_1^2 + O(\beta_1^3). \tag{5.2}$$

Recall that $\tau \in (1, \alpha_2)$. Thus, by (5.1) we have $|\tau - 1| = O(\beta_1)$. Combining this with (4.15), (5.2), and since by Corollary 4.1 and (3.7), we have that

$$\omega_{\beta_1,\beta_2} = n\tau\pi_1^*\omega_{\mathrm{FS}} + \varphi\Big(\pi_2^*\omega_{\mathrm{Cyl}} + \sqrt{-1}\alpha \wedge \overline{\alpha} + \sqrt{-1}\alpha \wedge \overline{dw/w} + \sqrt{-1}dw/w \wedge \overline{\alpha}\Big), \tag{5.3}$$

we conclude that

$$\lim_{\beta_1 \to 0} \omega_{\beta_1,\beta_2} = n\pi_1^*\omega_{\mathrm{FS}}, \tag{5.4}$$

and the convergence of tensors occurs smoothly, so $(\mathbb{F}_n \setminus (Z_n \cup Z_{-n}), \omega_{\beta_1,\beta_2})$ converges in the Gromov–Haussdorff sense to $(\mathbb{P}^1, n\omega_{\mathrm{FS}})$. (This combination of smooth convergence of tensors and simultaneous collapse in the Gromov–Haussdorff sense is reminiscent of [10, Theorem 2.4 (ii) (a)].) This concludes the proof of the first two statements in Theorem 1.2.

We now turn to the last statement in Theorem 1.2. Inspired by [12, Lemma 3.1], we change variable from $\tau \in (1, \alpha_2)$ to

$$y := \frac{\tau - 1 - \frac{n\beta_1}{2}}{n\beta_1^2/2},$$

with $y \in \left(-\frac{1}{\beta_1}, \frac{1}{\beta_1} + O(1)\right)$ (recall (5.1)), with $y = 0$ roughly corresponding to the mid-section between Z_n and Z_{-n}. Thus, by (4.15), (5.1), and (5.2), and since $|\tau - 1| = O(\beta_1)$,

$$\varphi(\tau) = \frac{2 - n\beta_1}{2n}\left(\frac{n^2\beta_1^2}{4} - \frac{n^2\beta_1^4}{4}y^2\right) + O(\beta_1^3), \quad \left(-\frac{1}{\beta_1}, \frac{1}{\beta_1} + O(1)\right). \tag{5.5}$$

To determine a fiberwise-rescaled limit, define the fiberwise-rescaled metric, where the rescaling only occurs for the terms that have a well-defined restriction to each fiber, i.e.,

$$\widetilde{\omega_{\beta_1,\beta_2}} := n\tau\pi_1^*\omega_{\mathrm{FS}} + \frac{1}{\beta_1^2}\varphi\pi_2^*\omega_{\mathrm{Cyl}} + \varphi\Big(\sqrt{-1}\alpha \wedge \overline{\alpha} + \sqrt{-1}\alpha \wedge \overline{dw/w} + \sqrt{-1}dw/w \wedge \overline{\alpha}\Big).$$

with φ given by (5.5). As in the proof of Lemma 3.2, the restriction of ω_{β_1,β_2} to a fiber is given by

$$\frac{1}{2\varphi(\tau)^2}d\tau \otimes d\tau + 2\varphi(\tau)^2 d\theta \otimes d\theta,$$

so the restriction of $\widetilde{\omega_{\beta_1,\beta_2}}$ to a fiber is given by

$$\frac{\beta_1^2}{2\varphi(y)^2}dy \otimes dy + 2\frac{\varphi(y)^2}{\beta_1^2}d\theta \otimes d\theta, \tag{5.6}$$

The Gromov–Haussdorff limit is then diffeomorphic to $\mathbb{P}^1 \times \mathbb{C}^\star$. Indeed, as in the proof of Lemma 3.2, and using (5.5), the length of the path on each fiber between the intersection point of the fiber with Z_n and its midpoint, the intersection of the fiber with the smooth section $\{y = 0\}$ is

$$\int_1^{1+n\beta_1/2} \sqrt{\varphi(\tau)}d\tau = \int_0^{n\beta_1/2} \sqrt{\varphi(y)}dy = O(\beta_1^2) \tag{5.7}$$

from Z_n, and, moreover, we claim these distance are comparable across fibers in the sense that $O(\beta_1^2)$ is uniform, i.e., bounded from above and below by some $C_1\beta_1^2$ and $C_2\beta_1^2$ with $C_1, C_2 > 0$ constants. To see this,

$$\int_0^{n\beta_1/2} \sqrt{\varphi(y)}dy = \sqrt{\frac{2 - n\beta_1}{2n}} \int_0^{n\beta_1/2} \sqrt{\left(\frac{n\beta_1}{2} - y\right)\left(\frac{n\beta_1}{2} + y\right)}dy + o(\beta_1^2). \tag{5.8}$$

Noting that

$$\int_0^{n\beta_1/2} \sqrt{\left(\frac{n\beta_1}{2} - y\right)}dy = \frac{2}{3}\left(\frac{n\beta_1}{2}\right)^{3/2},$$

we see that (5.8) is bounded from above by

$$\sqrt{\frac{1}{n}}\sqrt{n\beta_1} \int_0^{n\beta_1/2} \sqrt{\left(\frac{n\beta_1}{2} - y\right)}dy + o(\beta_1^2) = \frac{2}{3}\left(\frac{n}{2}\right)^{3/2}\beta_1^2 + o(\beta_1^2),$$

and from below by

$$\sqrt{\frac{1}{2n}}\sqrt{n\beta_1/2} \int_0^{n\beta_1/2} \sqrt{\left(\frac{n\beta_1}{2} - y\right)}dy + o(\beta_1^2) = \frac{1}{3}\left(\frac{n}{2}\right)^{3/2}\beta_1^2 + o(\beta_1^2),$$

as claimed. Thus, rescaling by $O(\beta_1^{-2})$ as we do in the definition of $\widetilde{\omega_{\beta_1,\beta_2}}$ we must get the product differential structure on $\mathbb{P}^1 \times \mathbb{C}^\star$ in the pointed Gromov–Hausdorff limit, as claimed.

Moreover, a computation similar to the one above shows that the distance between the mid-section and each of $Z_{\pm n}$ along a fiber is $O(\beta_1^{-1})$ (again, uniform across fibers) in the fiber-wise rescaled metric. Now, in the limit $\beta_1 \to 0$, (5.6) converges in the pointed (with the point being $y = 0$) sense on compact subsets to

$$\frac{2}{n}dy \otimes dy + \frac{n}{2}d\theta \otimes d\theta,$$

so changing coordinate to $t := y/(n/2)$ we obtain $\frac{n}{2}(dt \otimes dt + d\theta \otimes d\theta)$, whose Kähler form is $nd\zeta \wedge \overline{d\zeta}/|\zeta|^2 = n\omega_{\text{Cyl}}$ (recall (3.3)) with $\zeta := e^{t+\sqrt{-1}\theta}$. Combining this with (5.3), (5.5) and (5.6), the limiting metric, in the pointed Gromov–Hausdorff sense is then

$$\lim_{\beta_1 \to 0} \widetilde{\omega_{\beta_1,\beta_2}} = n(\pi_1^*\omega_{\text{FS}} + \pi_2^*\omega_{\text{Cyl}}).$$

This concludes the proof of Theorem 1.2.

Remark 5.1. We mention in passing an interesting borderline phenomenon that occurs in the cases $n \in \{1, 2\}$. By Corollary 4.1, when $n = 2$ the metrics $\omega_{\beta_1, \beta_2}$ exists for *all* $\beta_1 \in (0, 1)$. Naturally, one may ask about the *large angle limit* $\beta_1 \to 1$. It turns out that in this case $\alpha_2 \to \infty$ and one obtains Gromov–Haussdorff convergence to the complete Ricci flat Eguchi–Hanson metric on the non-compact space $-2H_{\mathbb{P}^1}$ as $C_2 = Z_{-n}$ gets pushed-off to infinity. Similarly, when $n = 1$, the metrics $\omega_{\beta_1, \beta_2}$ exists for *all* $\beta_1 \in (0, 2)$. In the limit $\beta_1 \to 2$ one has $\beta_2 \to 1$ and the limit is now a complete Ricci flat metric on $-H_{\mathbb{P}^1}$ with angle 4π along Z_1. We discuss these examples in detail elsewhere [13].

References

[1] E. Calabi, Métriques Kählériennes et fibrés holomorphes, Ann. scient. Éc. Norm. Sup. 12 (1979), 269–294.

[2] E. Calabi, Extremal metrics, in: Seminar on Differential Geometry (S.-T. Yau, Ed.), Princeton University Press, 1982, pp. 259–290.

[3] I.A. Cheltsov, private communication, 2015.

[4] I.A. Cheltsov, Y.A. Rubinstein, Asymptotically log Fano varieties, Adv. Math. 285 (2015), 1241–1300.

[5] I.A. Cheltsov, Y.A. Rubinstein, On flops and canonical metrics, Ann. Sc. Norm. Super. Pisa Cl. Sci. (5) 18 (2018), 283–311.

[6] K. Fujita, On log K-stability for asymptotically log Fano varieties, Ann. Fac. Sci. Toulouse Math. 25 (2016), 1013–1024.

[7] F. Hirzebruch, Über eine Klasse von einfach-zusammenhängenden komplexen Mannigfaltigkeiten, Math. Ann. 124 (1951), 77–86.

[8] A.D. Hwang, M.A. Singer, A momentum construction for circle-invariant Kähler metrics, Trans. Amer. Math. Soc. 354 (2002), 2285–2325.

[9] T. Jeffres, R. Mazzeo, Y.A. Rubinstein, Kähler–Einstein metrics with edge singularities, (with an appendix by C. Li and Y.A. Rubinstein), Annals of Math. 183 (2016), 95–176.

[10] A. Pulemotov, Y.A. Rubinstein, Ricci iteration on homogeneous spaces, Trans. Amer. Math. Soc. 371 (2019), 6257–6287.

[11] Y.A. Rubinstein, Smooth and singular Kähler–Einstein metrics, in: Geometric and Spectral Analysis (P. Albin et al., Eds.), Contemp. Math. 630, Amer. Math. Soc. and Centre de Recherches Mathématiques, 2014, 45–138.

[12] Y.A. Rubinstein, K. Zhang, Small angle limits of Hamilton's footballs, Bull. London Math. Soc. 52 (2020) 189–199.

[13] Y.A. Rubinstein, K. Zhang, Eguchi–Hanson–Stenzel metrics as limits of Kähler–Einstein edge metrics, preprint, in preparation.

[14] G. Tian, Kähler–Einstein metrics on algebraic manifolds, Lecture Notes in Math. 1646 (1996), 143–185.

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